

So far we've derived q_{atomic}

$$\epsilon_{\text{atom}} = \epsilon_{\text{translation}} + \epsilon_{\text{electronic}}$$

$$q_{\text{atom}} = q_{\text{trans}} * q_{\text{electronic}}$$

$$\epsilon_{\text{trans}} = \frac{h^2}{8ma^2} (n_x^2 + n_y^2 + n_z^2) \quad \leftarrow \text{particle in a box}$$

$$q_{\text{trans}} = \sum_{n_x=1}^{\infty} \sum_{n_y=1}^{\infty} \sum_{n_z=1}^{\infty} e^{-\frac{\beta h^2}{8ma^2} (n_x^2 + n_y^2 + n_z^2)}$$

$$\approx \left(\int_0^{\infty} e^{-\beta h^2 n^2 / 8ma^2} dn \right)^3$$

Big Box assumption

$$q_{\text{trans}}(V, T) = \left(\frac{2\pi m k_B T}{h^2} \right)^{3/2} V \quad \Rightarrow \quad \epsilon_{\text{trans}} = \frac{3}{2} k_B T$$

Also: $q_{\text{elect}} = g_1 + g_2 e^{-\beta(\epsilon_2 - \epsilon_1)} + g_3 e^{-\beta(\epsilon_3 - \epsilon_1)}$

↑
Degeneracies

‡ setting 0 of energy to ground electronic state

Now, for collective properties:

$$Q(N, V, T) = \frac{[q_{\text{trans}}(V, T) q_{\text{elect}}(T)]^N}{N!}$$

$$U = \langle E \rangle = - \left(\frac{\partial \ln Q}{\partial \beta} \right)$$

thermo dynamic
symbol for
internal energy

to do this, we need $\ln Q$:

$$\ln Q = N \left[\ln q_{\text{trans}} + \ln q_{\text{elect}} \right] - \ln N!$$

$$= N \left[\frac{3}{2} \ln \left(\frac{2\pi m}{\beta h^2} \right) + \ln V + \ln q_{\text{elect}} \right] - \ln N!$$

$$\left(\frac{\partial \ln Q}{\partial \beta} \right) = \frac{3N}{2} \left(\frac{\beta h^2}{2\pi m} \right) \left(\frac{2\pi m}{h^2} \right) \left(\frac{-1}{\beta^2} \right) + \frac{N \left(\frac{\partial q_{\text{elect}}}{\partial \beta} \right)}{q_{\text{elect}}}$$

$$U = \frac{3Nk_B T}{2} + \underbrace{N \left(g_2 e^{-\beta(\epsilon_2 - \epsilon_1)} + \dots \right)}_{q_{\text{elect}}}$$

← mostly very small

$$U \approx \frac{3}{2} RT$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_{N,V} = \frac{3}{2} R$$

One other interesting feature $\text{Pressure} = k_B T \left(\frac{\partial \ln Q}{\partial V} \right)_{N,T}$

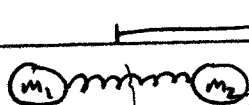
$$P = k_B T \frac{\partial}{\partial V} \left\{ \left[\frac{3}{2} \ln \left(\frac{2\pi m}{\beta h^2} \right) + \ln V + \ln q_{\text{elect}} \right] - \ln N! \right\}$$

$$P = k_B T \cdot \frac{1}{V}$$

← ideal gas law is a result of particle in a box energy levels!

Diatomic Molecules

4 kinds
of motion:

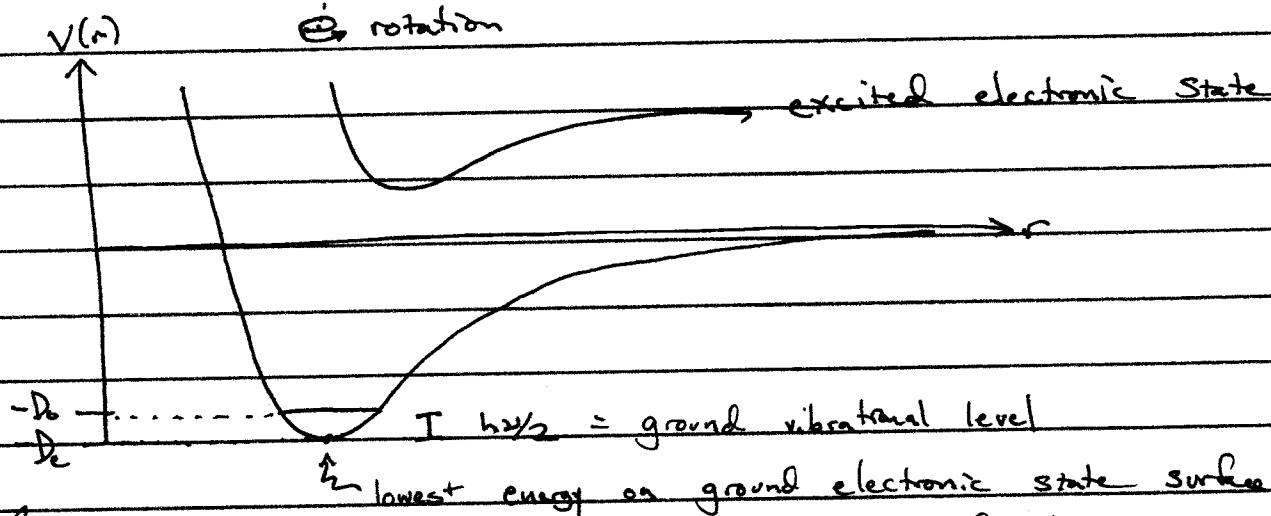


vibration

rotation

translation

+ electronic excitation



$\sum D_e = \text{dissociation energy from bottom of electronic state}$
 $D_0 = \text{dissociation energy from lowest vibrational state}$

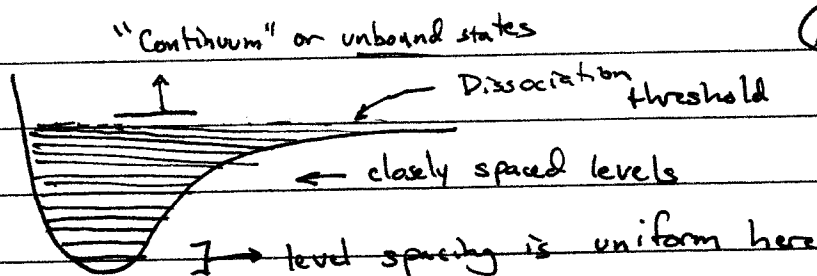
We'll leave: $q_{\text{elect}} = \sum_i g_i e^{-\beta \epsilon_i}$ $\epsilon_i = \text{energy of level } i$
 $g_i = \text{degeneracy of level } i$

Note that for molecules the 0 of energy is taken to be the dissociated atoms, not 0 for ground state

Translation is identical to atoms: $q_{\text{trans}} = \left(\frac{2\pi (m_1+m_2) k_B T}{h^2} \right)^{3/2} V$ ^{total mass}

Rotation & Translation are new.

Vibrational Motion



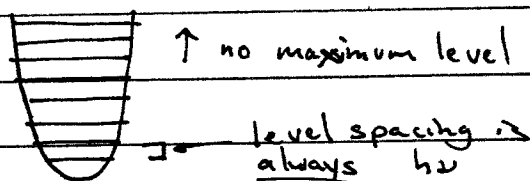
(10-4)

$$E_v = (v + \frac{1}{2}) h\nu - h\nu x_e (v + \frac{1}{2})^2 \quad \leftarrow \text{levels for a Morse oscillator}$$

The Harmonic model

$$E_v = (v + \frac{1}{2}) h\nu$$

$$v = 0, 1, 2, \dots, \infty$$



$$q_{vib} = \sum_{v=0}^{\infty} e^{-\beta h\nu (v + \frac{1}{2})} \quad \leftarrow \text{Boltzmann factor for that state}$$

↑ sum over states

$$= \sum_{v=0}^{\infty} e^{-\beta h\nu/2} e^{\beta h\nu v} \quad \leftarrow \text{first factor does not depend on } v$$

$$= e^{-\beta h\nu/2} \sum_{v=0}^{\infty} e^{-\beta h\nu v}$$

$$= \sqrt{X} \sum_{v=0}^{\infty} X^v$$

Let $X = e^{-\beta h\nu}$

← This is a geometric series!

$$= \sqrt{X} \frac{1}{1-X}$$

$$q_{vib} = \frac{e^{-\beta h\nu/2}}{1 - e^{-\beta h\nu}} = \frac{e^{-h\nu/2k_B T}}{1 - e^{-h\nu/k_B T}}$$

(10-5)

$\frac{h\nu}{k_B}$ has units of temperature

$$\Theta_{\text{vib}} = \frac{h\nu}{k_B}$$

$\nwarrow$ constant $\nwarrow$ property of molecule
 $\nearrow$ constant
 = "vibrational" temperature

$$q_{\text{vib}} = \frac{e^{-\Theta_{\text{vib}}/2T}}{1 - e^{-\Theta_{\text{vib}}/T}}$$

$\leftarrow$ functional form is independent of molecular species!

Also: $h\nu = \hbar\omega = \hbar\sqrt{\frac{k}{\mu}}$

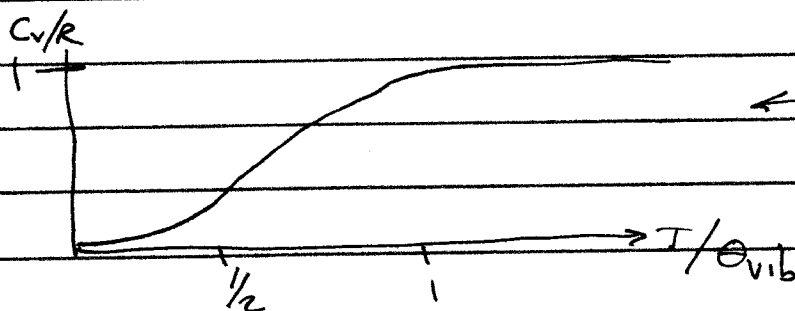
$\nwarrow$ spring constant
 $\nwarrow$ reduced mass

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

Average vibrational energy:

$$\begin{aligned} \langle E_{\text{vib}} \rangle &= -N \left(\frac{\partial \ln q_{\text{vib}}}{\partial \beta} \right) = N k_B T^2 \left(\frac{\partial \ln q_{\text{vib}}}{\partial T} \right) \\ &= N k_B \left(\frac{\Theta_{\text{vib}}}{2} + \frac{\Theta_{\text{vib}}}{e^{\Theta_{\text{vib}}/T} - 1} \right) \end{aligned}$$

$$C_V = \frac{\partial \langle E_{\text{vib}} \rangle}{\partial T} = R \left(\frac{\Theta_{\text{vib}}}{T} \right)^2 \frac{e^{-\Theta_{\text{vib}}/T}}{(1 - e^{-\Theta_{\text{vib}}/T})^2}$$



$\leftarrow$ vibrations can only hold heat when we exceed $\frac{\Theta_{\text{vib}}}{2}$