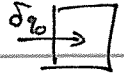


4 different processes can change an ideal gas
We will assume each one is performed reversibly.)

1) Constant Volume (isochoric)

initial state P_1, V_0, T_1 $\xrightarrow{\text{add heat}}$ δq with V_0 fixed 

Work can't be done because the volume is fixed.

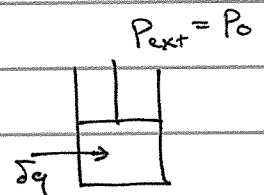
$$\delta w = 0 \quad \longleftarrow \quad dV = 0$$

$$\therefore dU = \delta q = C_v (T_2 - T_1) = \frac{C_v}{Nk} V_0 (P_2 - P_1)$$

So if we are given N, P_1, V_0, T_1 and δq , we can compute: $\delta w, P_2, T_2, \Delta U$, and $\Delta S = C_v \ln\left(\frac{T_2}{T_1}\right)$

2) Constant pressure (isobaric)

initial state P_0, V_1, T_1 $\xrightarrow{\text{add heat}}$ δq with a constant external pressure



Now work can be done:

$$\delta w = - \int_{V_1}^{V_2} P_0 dV = -P_0 (V_2 - V_1) \quad \text{Eq. (1)}$$

$\uparrow$ assumes a quasistatic process such that

$$P_{\text{int}} = \frac{NkT}{V} = P_0$$

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$$\Delta U = C_V (T_2 - T_1) = \frac{C_V P_0}{Nk} (V_2 - V_1)$$

Eq. (2)

$$\delta q = \Delta U - \delta w = \left(\frac{C_V}{Nk} + 1 \right) P_0 (V_2 - V_1)$$

Eq. (3)

So if we are given P_0, V_1, C_V, N , and δq

Eq. (3) above gives us V_2 , Eq. (1) gives us δw

Eq. (2) gives us T_2 and ΔU , and finally

$$\Delta S = C_V \ln \left(\frac{T_2}{T_1} \right)$$

3) Constant temperature (isothermal)

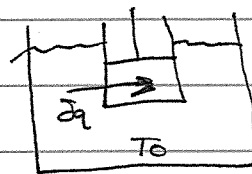
initial state

P, V, T_0



add heat

δq
with a constant
temperature



quasistatic implies: $P_{\text{ext}} = P_{\text{int}} = \frac{NkT_0}{V}$

$$\delta w = - \int_{V_1}^{V_2} P_{\text{ext}} dV = - \int_{V_1}^{V_2} P_{\text{int}} dV$$

$$= - \int_{V_1}^{V_2} \frac{NkT_0}{V} dV = -NkT_0 \ln \left(\frac{V_2}{V_1} \right)$$

ideal gas law

$$= -NkT_0 \ln \left(\frac{P_1}{P_2} \right)$$

$\Delta U = C_V \Delta T = 0$ for isothermal changes

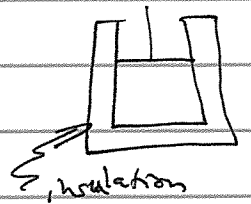
$$\therefore \delta q = -\delta w = NkT_0 \ln \left(\frac{V_2}{V_1} \right)$$

$$\Delta S = \frac{\delta q}{T_0} = Nk \ln\left(\frac{V_2}{V_1}\right)$$

Once again we can calculate $P_2, V_2, \Delta U, \delta w$ and ΔS from the initial state N, P, T_0, V_1 and δq

4) Adiabatic processes No heat is transferred

$$\begin{array}{ccc} \text{initial state} & \xrightarrow{\delta q = 0} & \text{final state} \\ P_1, V_1, T_1 & & P_2, V_2, T_2 \end{array}$$



We are allowed to reversibly change one variable P, V , or T

$$dU = C_v dT = \cancel{\delta q} - PdV$$

↑ because $\delta q = 0$ this must be work

For an ideal gas:

$$C_v dT = -\frac{Nk_B T}{V} dV$$

$$\int_{T_1}^{T_2} C_v dT = \int_{V_1}^{V_2} -\frac{Nk_B T}{V} dV$$

OR move the $\frac{1}{T}$ before integ

$$\int_{T_1}^{T_2} \frac{C_v}{T} dT = \int_{V_1}^{V_2} -\frac{Nk}{V} dV$$

$$C_v \ln \frac{T_2}{T_1} = -Nk \ln\left(\frac{V_2}{V_1}\right)$$

or:

$$\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)^{Nk/C_v}$$

T & V changes are related to each other in an adiabatic expansion

$$\Delta S = 0 \quad \text{because} \quad \delta q_{\text{rev}} = 0$$

Pathways are combinations of these processes

The four processes we have analyzed are limited because fixing one variable does not allow us to reach an arbitrary final state (P_2, V_2, T_2) from an arbitrary initial state (P_1, V_1, T_1)

We can combine 2 or more processes to form a pathway:

for example:

$$P_1, V_1, T_1 \xrightarrow{\text{isobaric}} P_1, V_2, T_x \xrightarrow{\text{isochoric}} P_2, V_2, T_2$$

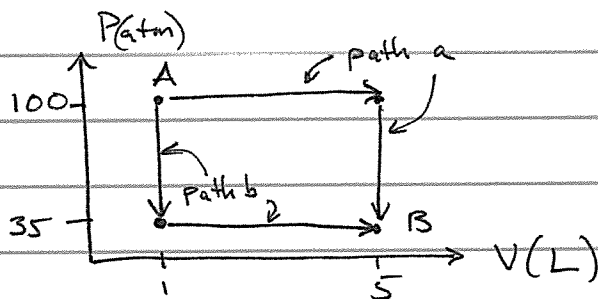
or

$$P_1, V_1, T_1 \xrightarrow{\text{isochoric}} P_2, V_1, T_y \xrightarrow{\text{isobaric}} P_2, V_2, T_2$$

These are 2 unique pathways that connect
 $P_1, V_1, T_1 \leftrightarrow P_2, V_2, T_2$

The values of state functions ΔU , ΔS will be identical for both paths, but δq & δw will be different.

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Example:

1 mol of ideal gas

state A: $P_1 = 100 \text{ atm}$, $V_1 = 1 \text{ L}$ state B: $P_2 = 35 \text{ atm}$, $V_2 = 5 \text{ L}$ path (a) isobaric $\rightarrow$ isochoric

$$\delta w_a = -P_1(V_2 - V_1) + 0 = -100(5 - 1) = -400 \text{ L} \cdot \text{atm}$$

path (b) isochoric $\rightarrow$ isobaric

$$\delta w_b = 0 - P_2(V_2 - V_1) = -(35)(5 - 1) = -140 \text{ L} \cdot \text{atm}$$

Work depended on the path

$$\text{state A: } T_1 = \frac{P_1 V_1}{R} \approx 1200 \text{ K}$$

$$\text{state B: } T_2 = \frac{P_2 V_2}{R} \approx 2100 \text{ K}$$

$$\begin{aligned} \therefore \Delta U &= C_V(T_2 - T_1) = \frac{3}{2}R(T_2 - T_1) = \frac{3}{2}R(2100 - 1200) \\ &= 111 \text{ L} \cdot \text{atm} \end{aligned}$$

This is the energy change regardless of path

$$\Delta U = \delta q + \delta w \quad \leftarrow \text{1st Law of thermodynamics}$$

Let's check this by computing δq directly:

$$\begin{aligned} \delta q &= C_V(T_2 - T_x) + P_1(V_2 - V_1) \leftarrow \text{isobaric} \\ &\quad + C_V(T_x - T_1) \leftarrow \text{isochoric} \end{aligned}$$

$$\delta q = C_V(T_2 - T_1) + P_1(V_2 - V_1)$$

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$$\delta q + \delta w = C_V (T_2 - T_1) + P_1 (V_2 - V_1) - P_1 (V_2 - V_1)$$

$$= C_V (T_2 - T_1) = \Delta U \quad \text{as expected}$$

Next time: Thermodynamic cycles